

# How to count ?

nodal domains of quantum graphs

Rami Band, Gregory Berkolaiko, Uzy Smilansky



מכון ויצמן למדע  
WEIZMANN INSTITUTE OF SCIENCE

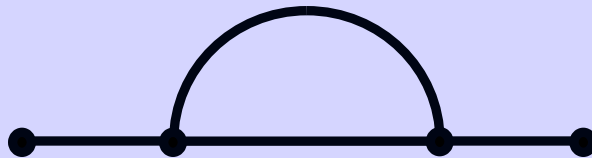
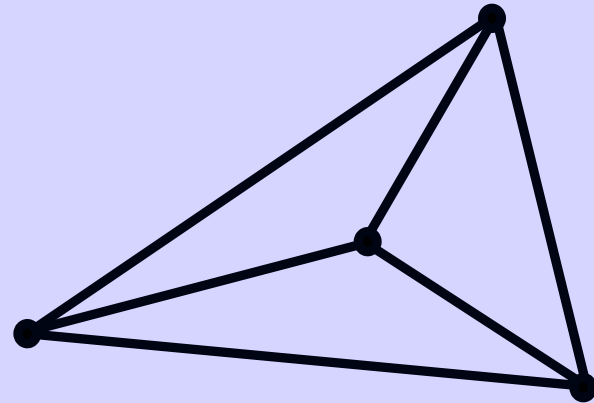
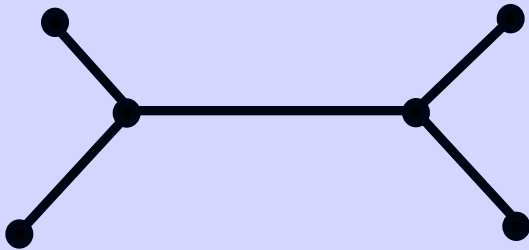


■ **1, 1, 2, 3, 5, 8, 13, 21, ?**

- What is the next number?
- Which formula generates this sequence?

■ **1, 2, 2, 4, 5, 5, 7, 7, 8, ?**

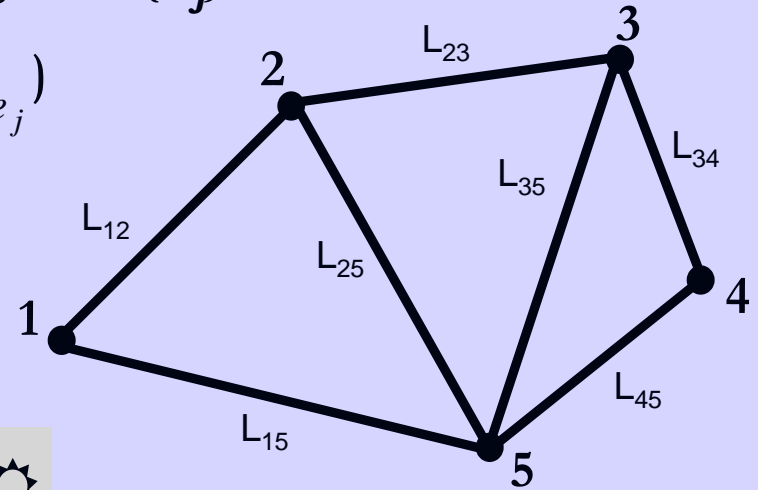
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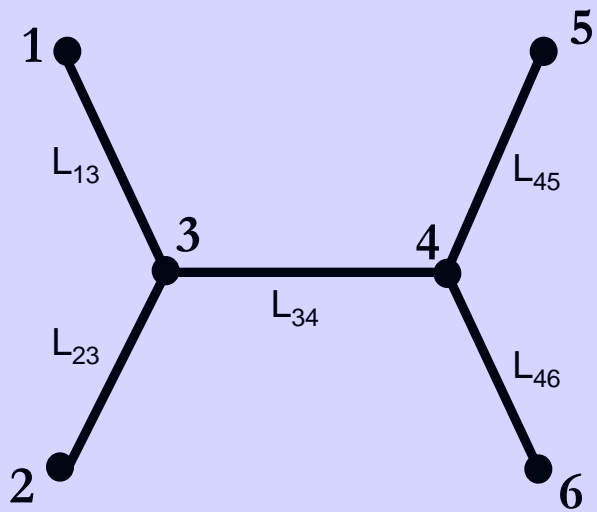
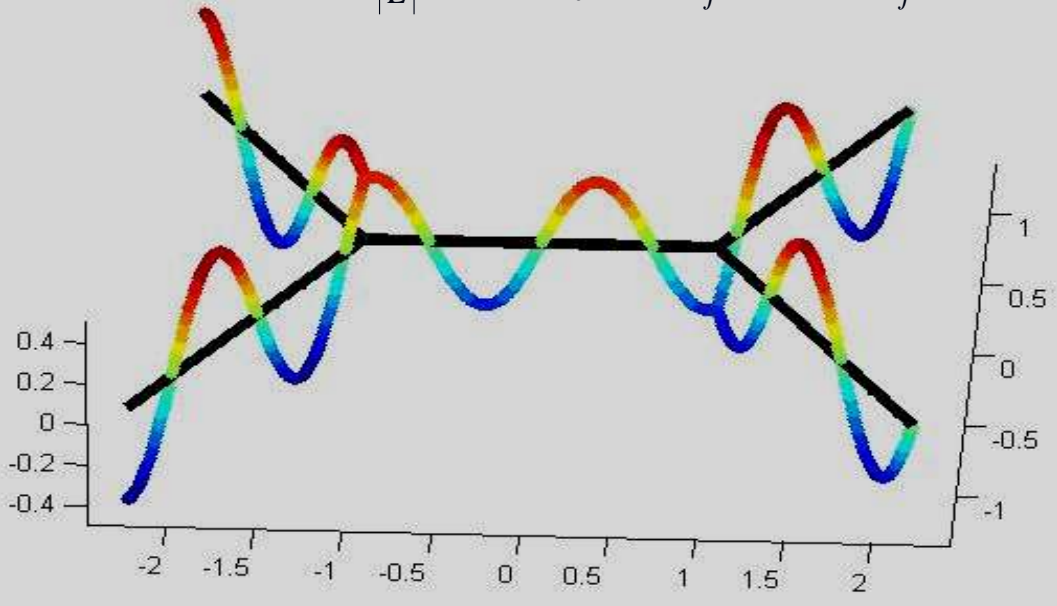
■ This sequence is the **nodal count**,  $\{v_n\}_{n=1}^{\infty}$ , of one of the graphs above.

# Metric Graphs - Introduction

- A **graph**  $\Gamma$  consists of a finite set of vertices  $V=\{v_j\}$  and a finite set of edges  $E=\{e_j\}$ .
- A **metric graph** has a finite length ( $L_{e_j}$ ) assigned to each edge.
- A **function** on the graph is a vector of functions on the edges:



$$f = (f_{e_1}, \dots, f_{e_{|E|}}) \quad \forall e_j, f_{e_j} : [0, L_{e_j}] \rightarrow \odot$$



# Quantum Graphs - Introduction

■ A **quantum graph** is a metric graph equipped with an operator, such as the negative **Laplacian**: 
$$-\Delta f = (-f''|_{e_1}, \dots, -f''|_{e_{|E|}})$$

■ For each vertex  $v$ , we impose **boundary conditions**, such as: (The **Neumann** boundary conditions)

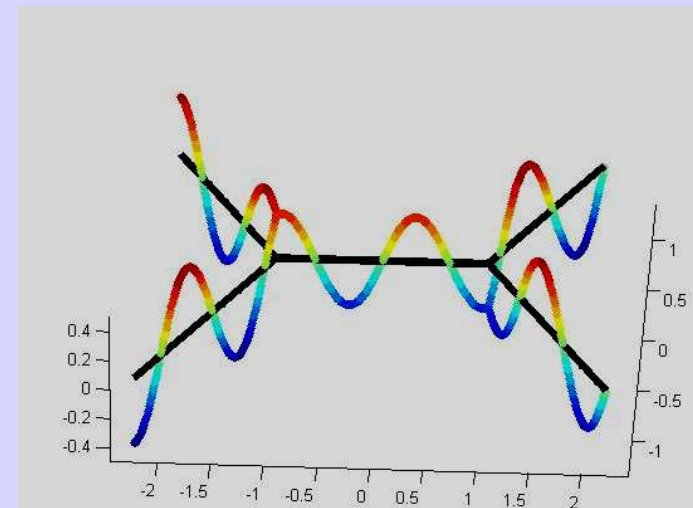
■ Continuity  $\forall e_1, e_2 \in E_v \quad f|_{e_1}(v) = f|_{e_2}(v)$

■ Zero sum of derivatives  $\sum_{e \in E_v} f'|_e(v) = 0$

■ For vertices of degree one, there are two special cases of boundary conditions, denoted

■ **Dirichlet**:  $f(v) = 0$

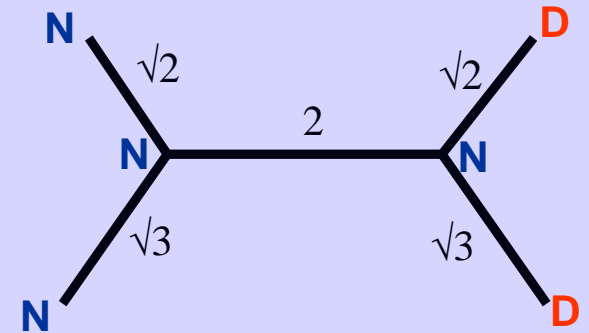
■ **Neumann**:  $f'(v) = 0$



# The eigenfunctions of Quantum Graphs

- A quantum graph is defined by specifying:
  - Metric graph
  - Operator
  - Boundary conditions for each vertex

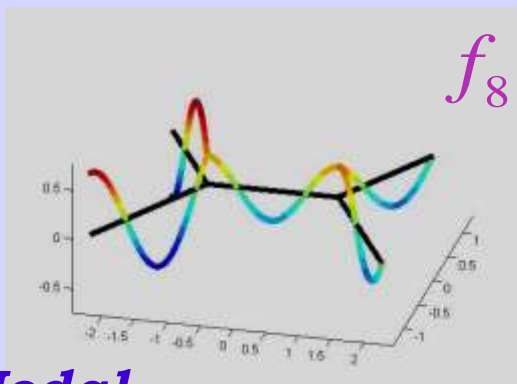
$$-\Delta = -\left( \frac{d^2}{dx^2} \Big|_{e_1}, \dots, \frac{d^2}{dx^2} \Big|_{e_{|E|}} \right)$$



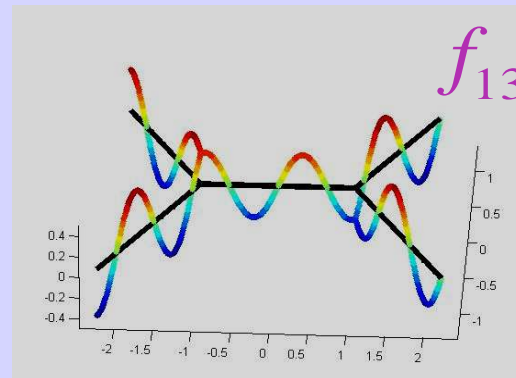
We are interested in the **eigenfunctions** of the Laplacian:

$$-\Delta f = k^2 f \Rightarrow \left( -f'' \Big|_{e_1}, \dots, -f'' \Big|_{e_{|E|}} \right) = \left( k^2 f \Big|_{e_1}, \dots, k^2 f \Big|_{e_{|E|}} \right)$$

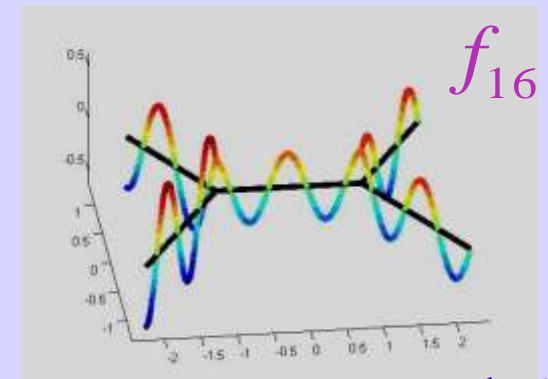
Examples of several eigenfunctions of the Laplacian on the graph above:



$f_8$



$f_{13}$



$f_{16}$

**Nodal count:**

$$\nu_8 = 8$$

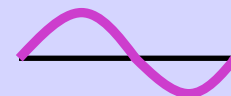
$$\nu_{13} = 13$$

$$\nu_{16} = 16$$

# The nodal count of Quantum Graphs

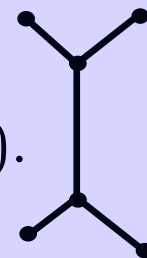
- We denote the nodal count sequence by  $\{v_n\}_{n=1}^{\infty}$ .

- The nodal count of a vibrating string is  $v_n = n$ .  
Sturm's oscillation theorem (1836).



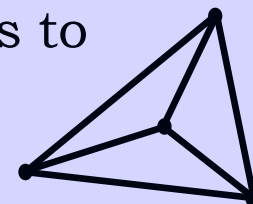
- A general bound of Courant (1923) is  $v_n \leq n$ .  
Adapted to quantum graphs by Gnutzmann, Smilansky, Weber (2004) following a method of Pleijel (1956).

- The nodal count of a tree graph is  $v_n = n$ .  
Al-Obeid, Pokornyi, Pryadiev (1992), Schapotschnikow (2006).



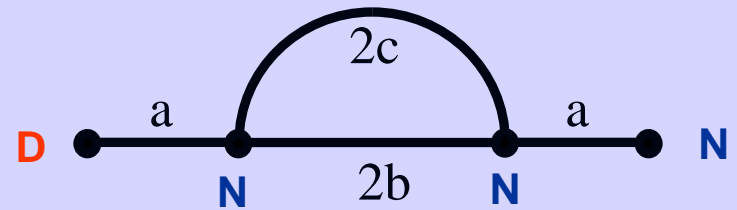
- $v_n \geq n - r$  is a bound given by Berkolaiko (2006),
  - where  $r = |E| - |V| + 1$  is the minimal number of edges to remove so that the graph turns into a tree.

$$r = 6 - 4 + 1 = 3$$



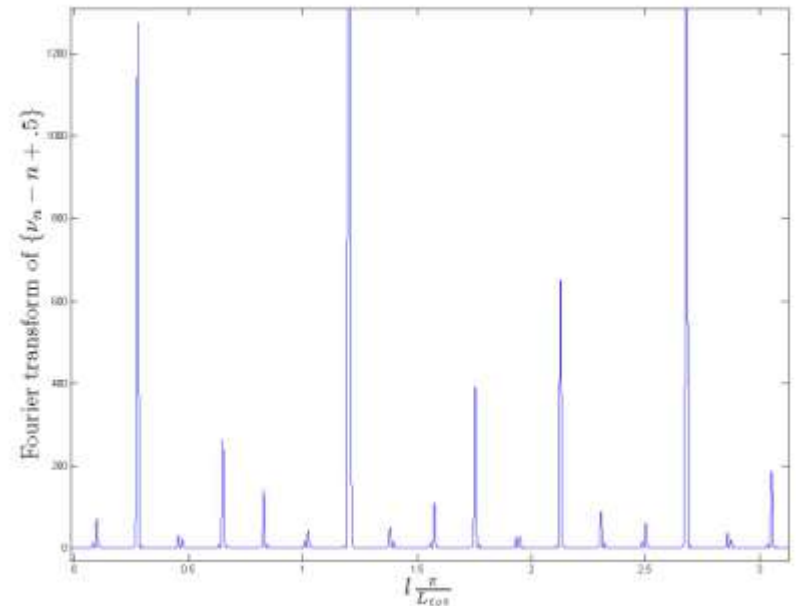
# The nodal count of Quantum Graphs

- Example - We try to deal with the nodal count of the graph



- The known bounds give  $n-1 \leq \nu_n \leq n$  ( $r=1$ ).
- We have the oscillating sequence  $\nu_n - n + 0.5 = \pm 0.5$ .
  - We calculate this sequence **numerically** (for some a,b,c values).
  - We Fourier transform it and obtain the following:

- Is there a formula for the nodal count of this graph ?



# The nodal count of a string

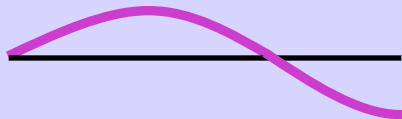
The eigenfunctions of

D ————— N

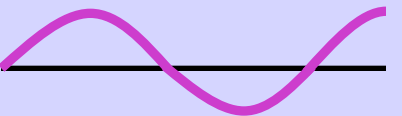
$k_1$



$k_2$



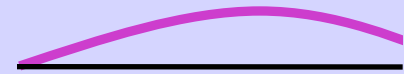
$k_3$



The eigenfunctions of

D ————— N

... Continuous  $k$  values ...

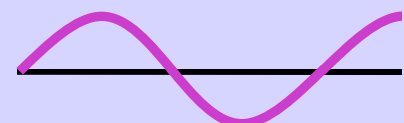
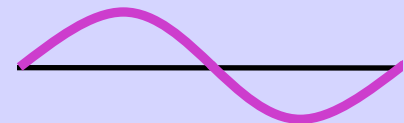
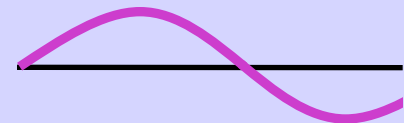
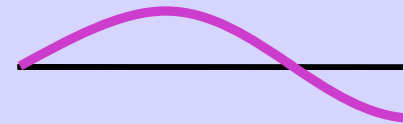
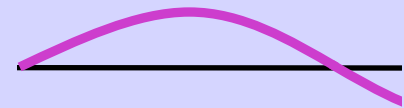
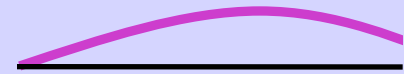




# Removing a boundary condition

The eigenfunctions of

**D** 

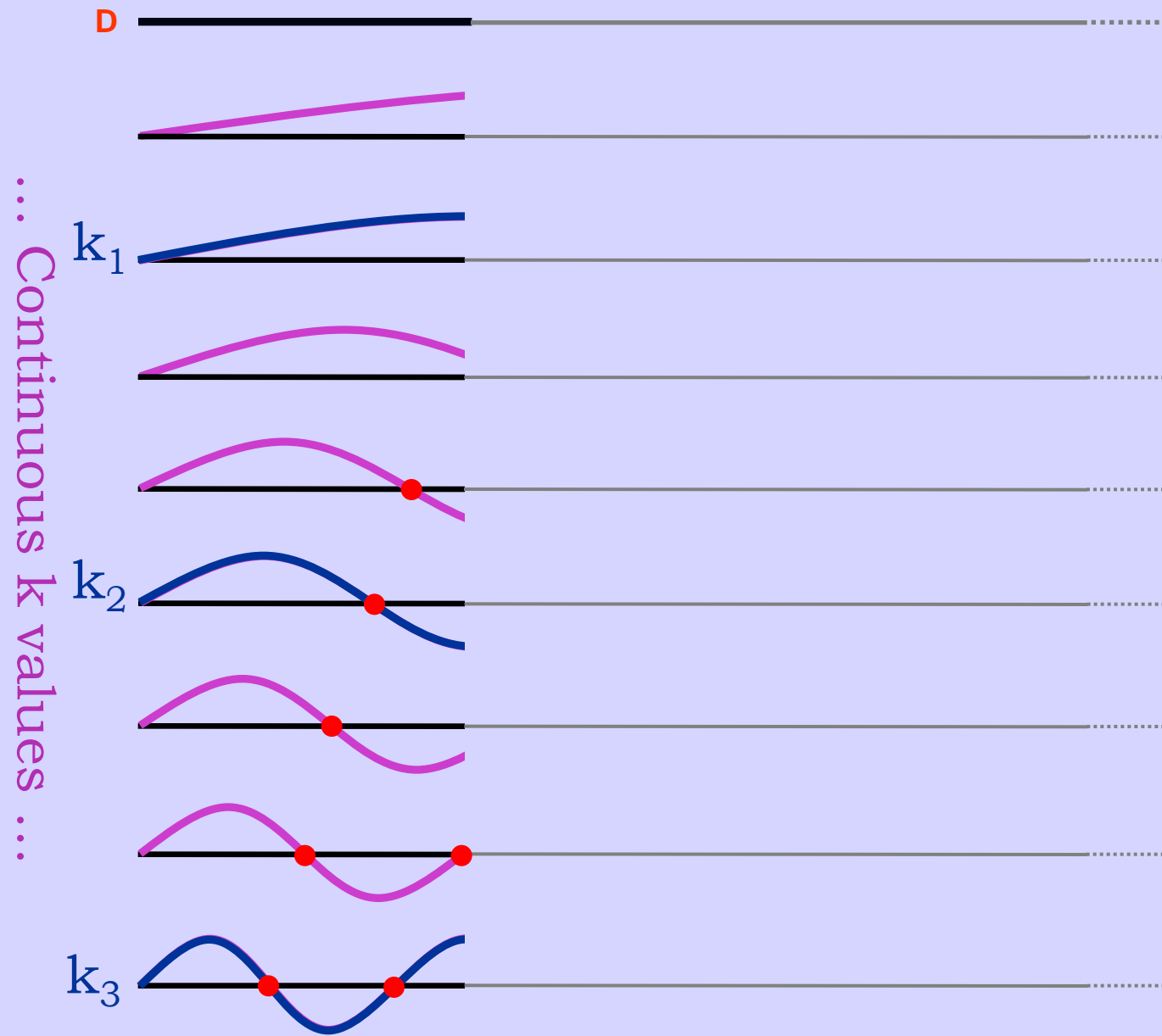


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# Removing a boundary condition

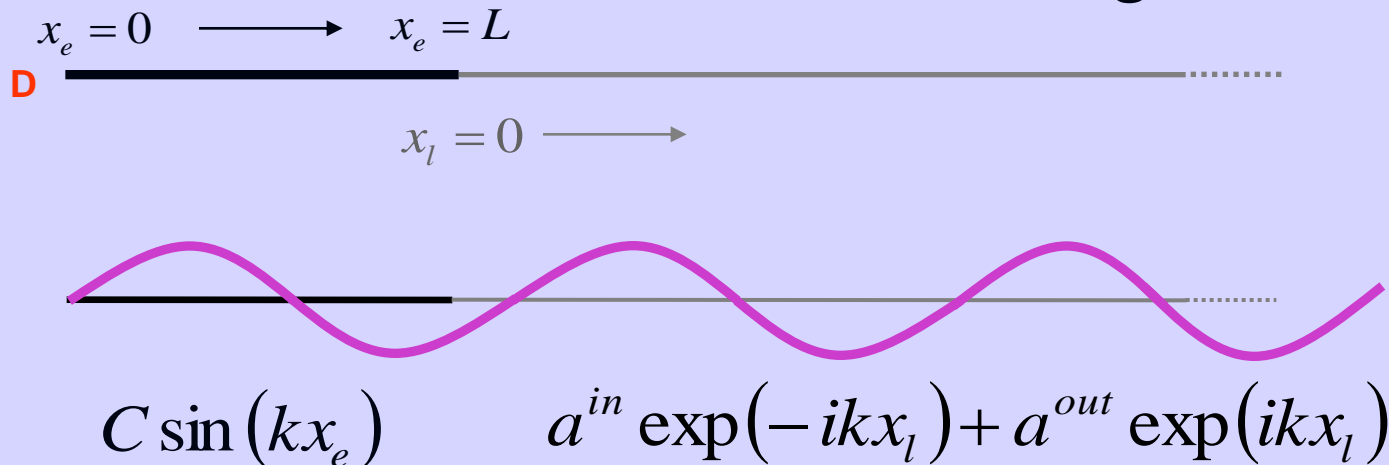
- attaching an infinite lead

The eigenfunctions of



# Removing a boundary condition

- attaching an infinite lead



$$\begin{cases} C \sin(kL) = a^{in} + a^{out} & \text{(Continuity)} \\ C \cos(kL) = ik(-a^{in} + a^{out}) & \text{(Zero sum of derivatives - differentiability)} \end{cases}$$

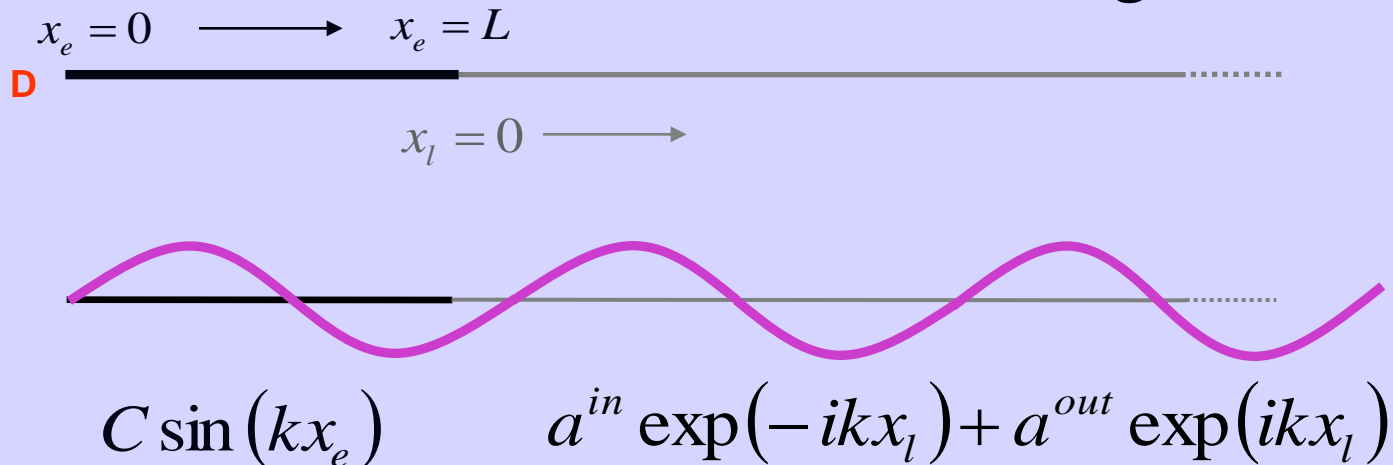
$$a^{out} = \exp(i(2kl + \pi)) a^{in} = S(k) a^{in} \quad S(k) = \exp(i\varphi(k)) \text{ is unitary.}$$

When  $\varphi(k) \equiv 0 \pmod{2\pi}$  we have an eigenfunction of the original graph.

When  $\varphi(k) \equiv \pi \pmod{2\pi}$  a nodal point enters the original graph.

# Removing a boundary condition

- attaching an infinite lead

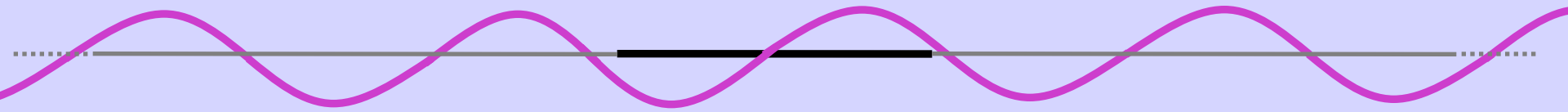


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# Removing two boundary conditions

- attaching two infinite leads

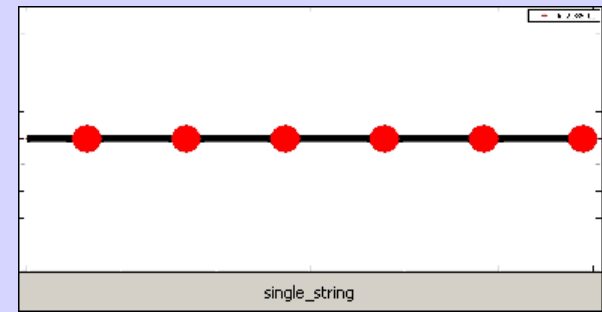


$$a_1^{in} \exp(-ikx_{l_1}) + a_1^{out} \exp(ikx_{l_1}) \\ = C_1 \sin(kx_{l_1} + \theta_1(k))$$

$$a_2^{in} \exp(-ikx_{l_2}) + a_2^{out} \exp(ikx_{l_2}) \\ = C_2 \sin(kx_{l_2} + \theta_2(k))$$

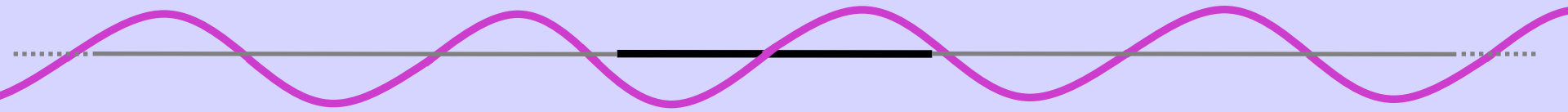
$$\begin{pmatrix} a_1^{out} \\ a_2^{out} \end{pmatrix} = S(k) \cdot \begin{pmatrix} a_1^{in} \\ a_2^{in} \end{pmatrix}$$

- For each  $k$  we have a two-dimensional space of possible functions.
  - In order to have a single valued movie we need to add a restriction.
  - We choose the restriction  $\theta_1(k) = \pi/2 + \theta_2(k)$  and get the movie:



# Removing two boundary conditions

- attaching two infinite leads

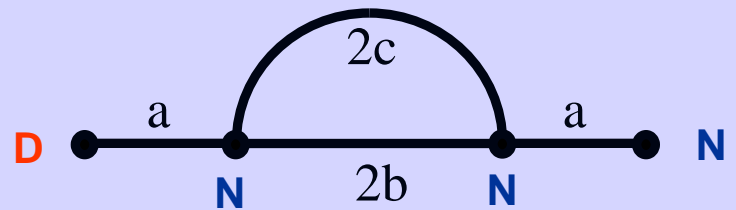


$$= C_1 \sin(kx_{l_1} + \theta_1(k))$$

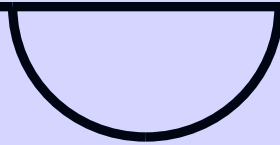
$$= C_2 \sin(kx_{l_2} + \theta_2(k))$$

- The interesting events are
  - For each  $k$  we have a two-dimensional space of possible functions.
    - In order to have a single valued movie we need to add a restriction.
    - We choose the restriction  $\theta_1(k) = \pi/2$ ;  $\theta_2(k) = 0 \pmod{\pi}$ 
      - Eigenfunction of the original graph & nodal point enters from the left.
      - Eigenfunction of the original graph & nodal point enters from the right.
- We can check that  $\theta_1'(k) > 0$  and  $\theta_2'(k) > 0$  and deduce that  $\mathcal{V}_n = n$ .

# Back to the nodal count of

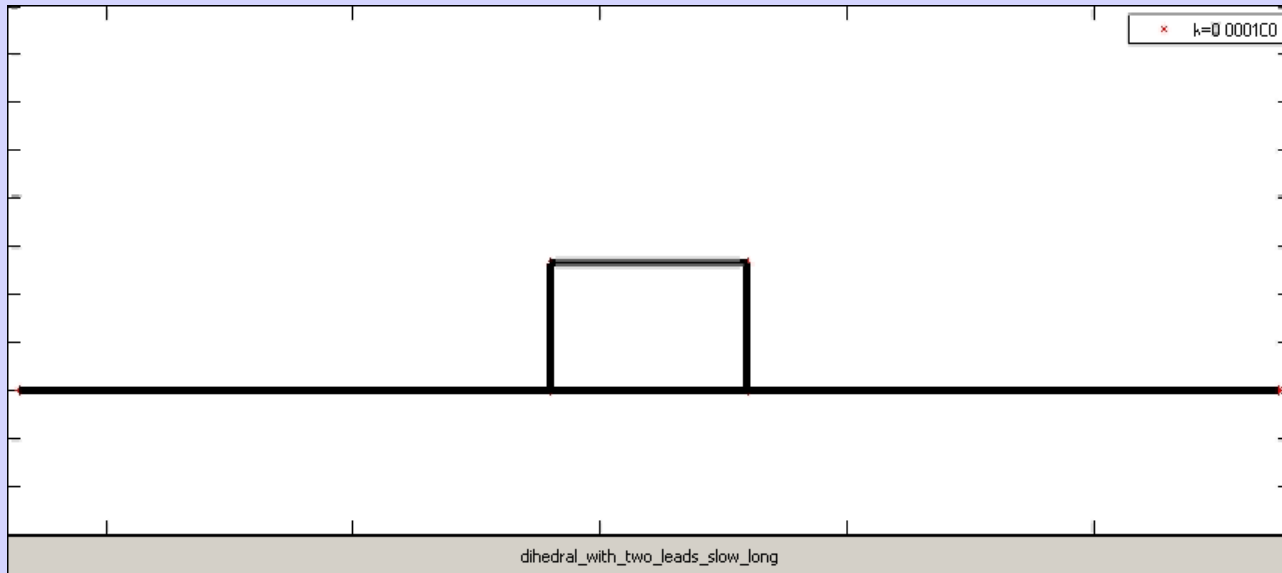


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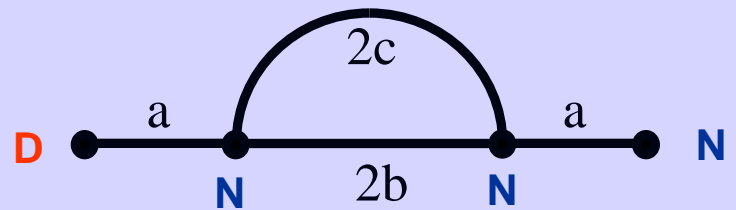
$$C_2 \sin(kx_{l_2} + \theta_2(k))$$

- The same restriction,  $\theta_1(k) = \pi/2 + \theta_2(k)$ , gives the following movie:

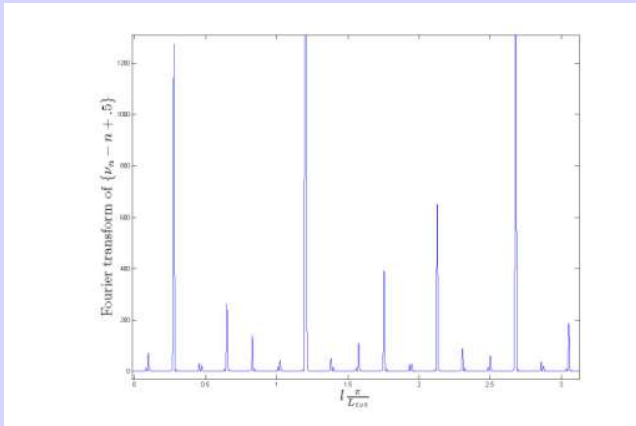


- Conclusions from the movie:
  - Nodal points enter the graph from the leads (as before).
  - An eigenvalue occur when a nodal point enters the graph (as before).
  - However, this time we also have splits and merges of nodal points.

# Back to the nodal count of



Analyzing the movie gives the formula:



$$v_n = n - 0.5 - 0.5 \cdot (-1)^{\left\lfloor \frac{b+c}{a+b+c} n \right\rfloor}$$

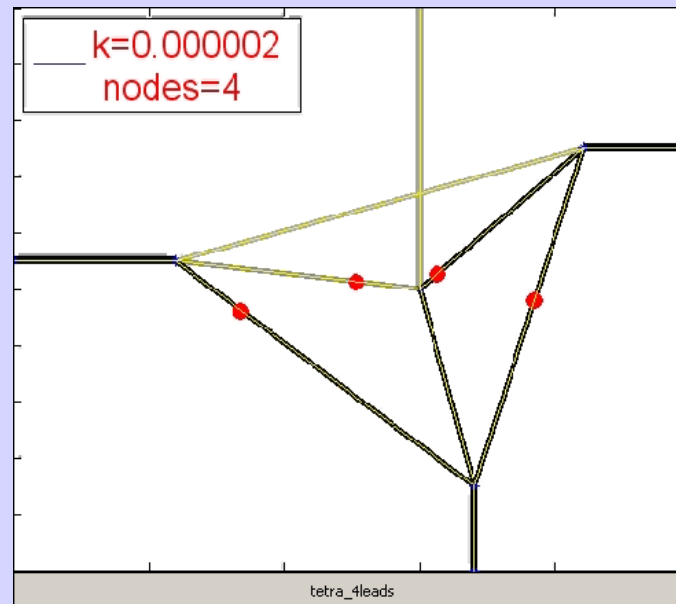
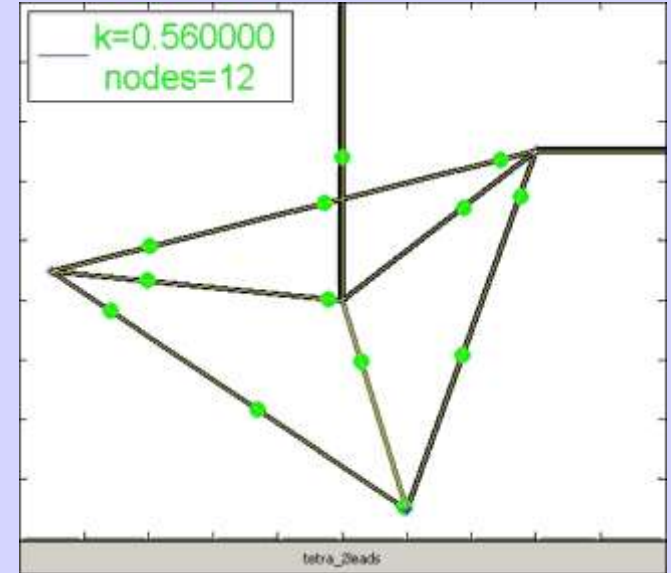
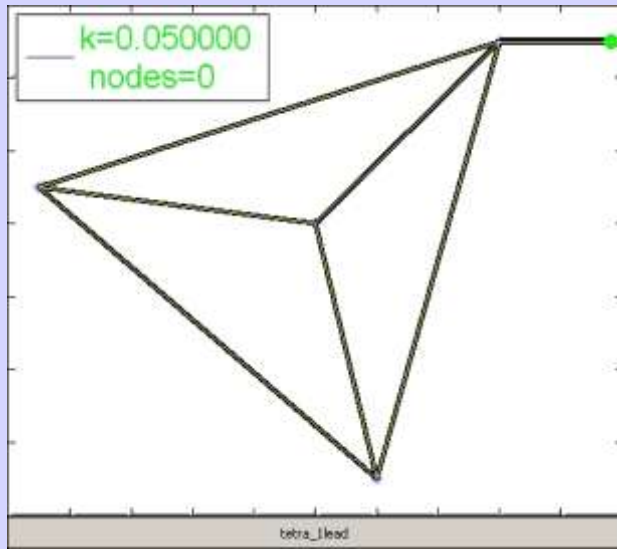
The nice Fourier transform is due to  $(-1)^{\lfloor x \rfloor} = \sum_{k=0}^{\infty} \frac{4}{(2k+1)\pi} \sin((2k+1)\pi x)$

which gives

$$v_n = n - 0.5 - \sum_{k=0}^{\infty} \frac{2}{(2k+1)\pi} \sin\left( (2k+1) \frac{b+c}{a+b+c} \pi n \right)$$



# What's next ?

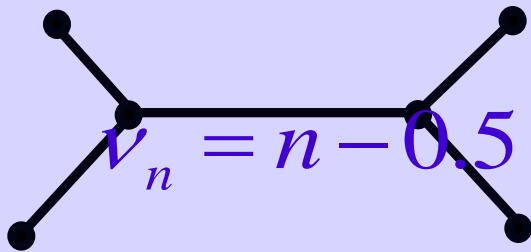


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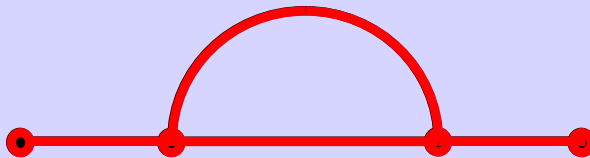
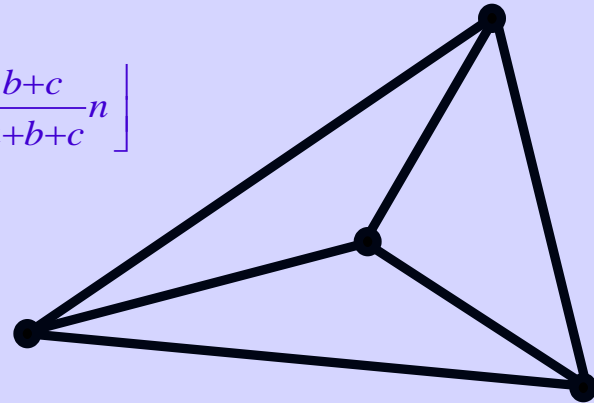
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$$v_n = n - 0.5 - 0.5 \cdot (-1)^{\lfloor \frac{b+c}{a+b+c} n \rfloor}$$



■ This sequence is the **nodal count**,  $\{v_n\}_{n=1}^{\infty}$ , of one of the graphs above.

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